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MAGNETIC ETCHING PROCESS, ESPECIALLY FORMAGNETIC OR MAGNETOOPTIC RECORDINGBACKGROUNDFIELD OF THE INVENTION

The present invention relates to a magnetic etching process.

More particularly, the invention applies advantageously to ultrahigh-density magnetic recording (production of discrete magnetic materials, magnetic memory circuits, magnetically-controllable logic circuits, etc.), optical recording of the read-only memory type (CDROM, DVDROM, etc.) and production of magnetically-controllable optical circuits (diffraction gratings, photonic gap materials, etc.) using a controlled variation of the optical index component associated with the magnetism.

PRIOR ART

The extraordinary development of multimedia technologies and services in recent years has led to a race to increase the recording density. In the field of rewritable disks, although optical (phase change) technologies are developing rapidly, magnetic techniques remain the first choice, and most particularly the "hard disk", for its high transfer rate. However, the magnetic techniques ought to be limited to storage densities of 100 bits/cm<sup>2</sup>.

One of the limiting factors will especially be the transition to contact recording, for distances between the read head and the recording medium of less than 10 nm: there is a trend toward recording technologies of "the" "tunnel-effect microscopy" ("STM-like storage") or "near-field" type.

Several technological jumps have been proposed in this direction in recent years, for example near-field CD-ROM or near-field magneto-optic recording.

In this regard, reference may advantageously be made to the following various publications:

Y. Martin, S. Rishton, H.K. Wickramasinghe, Appl. Phys. Lett. **71**, 1 (1997).

5 Y. Betzig, J.K. Trautman, T.D. Harris, J.S. Weiner, R.L. Kostelak, Science **251**, 1468 (1991).

B.D. Terris, H.J. Mamin, D. Rugar, W.R. Studenmund, G.S. Kino, Appl. Phys. Lett. **65**, 388 (1994).

10 E. Betzig et al., Appl. Phys. Lett. **61**, 142 (1992).

M. Myamoto, J. Ushiyama, S. Hosaka, R. Imura, J. Magn. Soc. Jpn. **19-S1**, 141 (1994).

T.J. Silva, S. Schultz, D. Weller, Appl. Phys. Lett. **65**, 658 (1994).

15 M.W.J. Prinz, R.H.M. Groeneveld, D.L. Abraham, H. van Kempen, H.W. van Kesteren, Applied. Phys. Lett. **66**, 1141 (1995).

Reference may also be made to the publication:

20 B.D. Terris H.J. Mamin, D. Rugar, Appl. Phys. Lett. **68**, 141 (1996) in which it was announced that the company 3M would shortly be commercializing a magnetooptically-read "hard disk" using a solid immersion lens (SIL).

25 However, the main limitation of magnetic techniques should be the "paramagnetic limit", that is to say the size below which the bits will be erased by themselves due to a thermal effect.

In the current hard disk technology, the

30 recording medium is a particulate material (magnetic particles in a nonmagnetic matrix, or magnetic particles (grains) separated by nonmagnetic grain boundaries (ME tape)). Now, minimization of the noise necessitates increasing the number of magnetic

35 particles seen by the read head, while these particles must be magnetically decoupled as far as possible. The size of the particles is therefore very much less than the size of a bit. By extrapolating the current data,

the particles would become paramagnetic below 8 nm, thereby limiting the recording density to around 100 bits/ $\mu\text{m}^2$ .

In magnetooptic recording, the materials used 5 at the present time are amorphous alloys of the rare earth/transition metal type, which could be replaced with Co/Pt multilayers or alloys with the advent of the blue laser. Bits 60 nm in size could actually be written by a thermomagnetic effect in continuous Co/Pt 10 multilayers, but it is probable that noise problems due to the recording medium (domain stability, domain wall roughness) would intervene, at bit sizes very much greater than 60 nm.

To extend this limit, it has recently been 15 proposed to replace the current recording medium materials with discrete materials in which the magnetic bit limits would be geometrically defined by lithographic methods:

either deposition on an etched surface,  
20 S. Gadetsky, J.K. Erwin, M. Mansuripur,  
J. Appl. Phys. 79, 5687 (1996)

or growth of isolated magnetic particles whose size and position are defined by lithography,

25 S.Y. Chou, M.S. Wei, P.R. Krauss, P. Fischer,  
J. Appl. Phys. 76, 6673 (1994).

The latter technique would allow there to be only a single magnetic particle per bit.

In parallel, pressing techniques based on a matrix defined by electronic lithography have been 30 developed,

S.Y. Chou, P.R. Krauss, P.J. Renstrom, Science 272, 85 (1996),

Y. Xia, X.M. Zhao, G.M. Whitesides,  
Microelecton. Eng. 32, 255 (1996),

35 which, just as in X-ray or interferential lithography, could in the near future allow mass production of etched media, with patterns very much less than one micron in size over areas of a few  $\text{cm}^2$ , probably sufficient for disks of the future.

However, in the current published work, these various techniques have several drawbacks:

1. Whatever the technique adopted, recording in contact mode will require a material having a low and controlled surface roughness: the etched materials proposed up until now will therefore require a final, and probably difficult, planarization step.
2. In the case of near-field magnetooptic recording, sudden variations in optical index (variations in reflectivity) of the etched material will give diffraction effects, which may be manifested by much greater polarization variations than those induced by the magnetic domains - a source of unacceptable noise.
3. A final problem, at very high densities on these etched materials, concerns the following of the track, and it will probably be necessary to develop a specialized "track" for this purpose, but without degrading the points mentioned above.

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#### PRESENTATION OF THE INVENTION

The subject of the invention is a magnetic etching process, characterized in that a thin-film magnetic material (comprising a few atomic planes) is controllably irradiated in order to locally modify, over regions having a width of the order of one micrometer or less, the magnetic properties of said material, such as, in particular, its coercivity, its magnetic anisotropy or its Curie temperature.

Such a process allows the aforementioned problems to be solved. In particular:

1. The roughness of the original film is unchanged by irradiation and can therefore be adjusted independently. In particular, it may be envisaged to carry out a postirradiation deposition (for the production of devices) under excellent growth conditions (% at an etched surface).

2. The optical index variations remain small for considerable changes in the magnetic properties and can, moreover, be controlled, within a certain range, almost independently of the magnetic variations 5 obtained, by the structure of the substrate or the energy of the ions.

3. The effect of the irradiation is cumulative: it is possible to carry out the irradiation several times, and to obtain the same result as in a 10 single time with the cumulative dose. This aspect may be useful when it is desired to irradiate several regions of the specimen with different values, or at different steps in the fabrication of a device.

4. The effect of the irradiation may be easily 15 controlled in real time, by measuring the change in the properties (for example magnetic properties) over a test region.

5. The technique is easy to employ for the mass production of recording media, and to do so 20 economically since the tools that it requires to be used are either already used in microelectronics (irradiation) or are under development (lithography by pressing in the case of large areas and of nanometric sizes, for example).

25 Advantageously, the irradiation is carried out by means of an ion beam.

Other technical means of energy deposition could be envisaged.

The irradiation may be carried out through a 30 resin mask or with the aid of a focused ion beam.

The aforementioned etching process is advantageously used for the ultrahigh-density magnetic or magnetooptic recording of binary information, and especially for the production of discrete magnetic 35 materials, of magnetic memory circuits or of magnetically-controllable logic circuits.

In particular, the aforementioned process has the advantage of making it possible to write magnetic domains of size very much less than 100 nm and whose

position and geometry are perfectly defined and therefore to maximize the signal-to-noise ratio and optimize the track-following problems, while preserving perfectly controlled surface roughness.

5 In addition, the process proposed by the invention is advantageously used for producing an optical recording of the read-only memory type (CDROM, DVDROM, etc.).

10 It is known in fact that the near-field optical recording techniques will probably have to use smooth writing materials, with a read head flying a few nm above said material (at the present time, 30 nm for a hard disk). Now, the current optical recording techniques of the read-only memory type are not 15 satisfactory: the pressing methods, using dies, may give sizes of less than 100 nm but the recording medium which is obtained is rough; as regards the writing methods using a focused laser beam (ablation, phase change), these do not make it possible to work with bit 20 sizes of the order of or less than 100 nm.

25 Applications other than the recording of binary information may be envisaged. In particular, the magnetic etching process proposed by the invention is advantageously used for the production of magnetically-controllable optical circuits (diffraction gratings, photonic gap materials, etc.) using a controlled variation of the optical index component associated with the magnetism, for the production of sensors (hard disk read heads, etc.) or magnetic memory circuits 30 (extraordinary Hall-effect memory, magnetoresistive memory, spin-dependent tunnel-effect memory).

35 In particular, it is known that the emergence of photonic gap materials opens the way to producing optical devices and that one of the aspects to be resolved will be that of control of the device. The process proposed by the invention makes it possible, by irradiation through a mask, to manufacture a waveguide film made of nonmagnetic material, comprising a regular array of magnetic units (photonic crystal) having an

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optical index which is both slightly different from that of the host material and magnetically controllable.

In general, the process proposed by the invention may apply whenever it is advantageous to define a magnetic element accurately, while maintaining a very high degree of planarity of the device (for example, in order to favor subsequent growth).

The process proposed by the invention may also be used for magnetically etching a layer already buried beneath other, insensitive layers, by adjusting the irradiation conditions. For example, and by way of nonlimiting indication, it is possible to produce electrical circuits etched in the same thin-film magnetic material, and only the important part of which will remain magnetic, the contact tracks having been made inactive by irradiation; the coercive field of a given region of a specimen may be controllably reduced so as to guarantee that the reversal of the magnetization will always occur under the same conditions, from the same site.

The process proposed by the invention may a priori be adapted to any material for which a minute variation in the local atomic arrangement can lead to a large modification in the magnetic properties, that is to say to transition metal alloys (e.g.: CoPt, NiFe, etc.), to rare earth/transition metal alloys (e.g.: TbFeCo, etc.) and to magnetic multilayers (e.g.: Co/Pt, Fe/Tb, etc.), without this list being exhaustive.

30 Co/Pt multilayers are materials which are potentially of interest for short-wavelength magnetooptic recording in blue light.

## DESCRIPTION OF ONE OR MORE EMBODIMENTS

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The process of magnetic etching by irradiation is described below in the case of magnetic multilayers irradiated by an ion beam and involves several steps, in which:

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- (i) the composition and the roughness at the interfaces and on the surface of the layers are carefully controlled before irradiation;
- 5 - (ii) the multilayer structure is irradiated by an ion beam, the structural modification induced by the beam being controlled; in particular, the energy density deposited by the beam is controlled by choosing the mass and the energy of the incident ions;
- 10 - (iii) the irradiation may be completed by a suitable thermal annealing step in order to relax the stresses and/or induce local ordering.

In the case of magnetic materials, the effects of the process are important on alloys (transition metal alloys, rare earth alloys and rare earth/transition metal alloys) and on multilayer stacks of all types.

20 The process is advantageously employed on Co/Pt multilayers. It should be noted that these materials have already been very widely studied for their properties, firstly their perpendicular magnetic anisotropy and secondly their strong magnetooptic Kerr effect; they therefore constitute advantageous candidates for magnetooptic recording.

25 In materials based on ultrathin multilayer films, the properties are dominated by the competition between the interface effects and the volume properties. For example, the easy magnetization direction is given by the sign of an effective anisotropy coefficient  $K_{eff}$  which, to a first 30 approximation, is given by:

$$K_{eff} = -K_d + K_v + \frac{(K_{s1} + K_{s2})}{t_{co}}$$

35 The first term represents the dipole shape anisotropy ( $K_d > 0$ ), the second term represents the volume anisotropy ( $K_v > 0$  in the case of Co) and the last term is due to the interfaces ( $K_s > 0$  in the case of the Co/Pt interface), the influence of which varies

inversely with the Co thickness  $t_{Co}$  ( $K_{s1}$  and  $K_{s2}$  denoting the magnetic anisotropy coefficients of the two interfaces of the Co film. Depending on the sign of  $K_{eff}$ , the easy magnetization axis is either the axis perpendicular to the plane of the layers ( $K_{eff} > 0$ ) or the plane of the film. The perpendicular configuration is necessary for magnetooptic recording and will probably become the standard for ultrahigh-density magnetic recording, all techniques included.

The process is preferably limited to irradiation resulting in low energy deposition (small number of atomic displacements at the interfaces that we are interested in). This may be achieved, for example, by light ions (e.g.  $He^+$ ) of low energy (from a few keV to about a hundred keV) or else by heavy ions (e.g. mass of the order of 100) of relatively high energy (typically, 1 MeV). The irradiation firstly modifies the composition of the interface and therefore, in particular, the anisotropy. For the thinnest films (1 or 2 atomic planes) or for higher doses, the composition of the film and hence its volume magnetism are also modified (by transferring atoms from one layer to another): in the particular case of Co/Pt, the Curie temperature of the CoPt alloy decreases with Pt concentration, and becomes below room temperature at around 75% Pt.

For example, the inventors have rendered specimens, having a thickness  $t_{Co}$  of 0.5 nm, paramagnetic at ordinary temperature, in a controlled manner, by irradiating, at a (very low) dose of  $10^{15}$  ions/cm<sup>2</sup>, with  $Kr^-$  ions accelerated to 300 keV, as well as with 30 keV  $He^+$  ions at a dose of  $10^{16}$  ions/cm<sup>2</sup>.

The effects of the irradiation were firstly characterized on simple Pt(3.4 nm)/Co( $t_{Co}$ )/Pt(6.5 nm)/amorphous substrate (Herasil polished silica,  $SiO_2/Si$ ,  $Si_3N_4/Si$ ) sandwiches deposited by sputtering.

With the deposition technique used, magnetic films with a perpendicular easy magnetization axis and a perfectly square polar hysteresis loop (100% remanent

magnetization) within the Co thickness range: 0.3 - 1.2 nm are obtained before irradiation.

The irradiation of these specimens at  $\text{He}^+$  ion fluences up to around  $2 \times 10^{15}$  atoms/cm<sup>2</sup>, the ions being 5 accelerated to energies of between 5 and 100 keV, makes it possible actually to adjust the magnetic properties of an ultrathin Co layer:

1. on 0.5 nm thick layers (approximately 10 2.25 atomic planes), the main effect is a drop in the Curie temperature, which may fall below room 15 temperature for a dose of the order of  $2 \times 10^{16}$  ions per cm<sup>2</sup>. Below that, the film retains a perpendicular easy magnetization axis and a square loop, but the coercive field of which decreases uniformly when the irradiation 20 dose is increased. Square magnetization loops with coercivities of a few Oe have been obtained. Advantageous applications for the production of low-field sensors may be envisaged;

2. on 1 nm thick specimens (approximately 25 5 atomic planes), the main effect of the irradiation is a tilt of the easy magnetization axis in the plane of the film, combined with a reduction in the interface anisotropy term  $K_s$ . The effect is obtained for low doses because the initial thickness is close to that 30 (1.2 nm) at which the tilting effect occurs in the original specimens;

3. on specimens of intermediate thickness (0.8 nm, i.e. 4 atomic planes), the same doses have no 35 visible effect on the hysteresis loop: at these thicknesses, the Curie temperature is already very high (close to that of bulk Co), and therefore largely insensitive to small modifications of the interface, these thicknesses also being very far from the natural thickness for tilting of the easy magnetization axis. This constitutes a useful characteristic of the process 40 since it makes it possible, on the one hand, to irradiate a bilayer while modifying only one of the layers and, on the other hand, to work at much higher doses, more conducive to homogeneity.

It should be noted that the acceleration energy of the ions has a lesser effect on the modification of the magnetic properties than on the depthwise distribution of the level of displacements in the material. This may allow the process to be employed in thin layers buried at substantially greater depths than those used in the demonstration example.

An essential characteristic of the process proposed is that, although the effect of the irradiation on the magnetism is great, its effect on the optical reflectivity of the specimen remains small.

The contrast is invisible to the naked eye, and barely visible in a good microscope (contrast comparable to that of a domain wall in a Pt/Co/Pt specimen). The smallness of the optical effect is due to the smallness of the induced structural modifications.

Tests on  $(\text{Pt}/\text{Co})_6/\text{Pt}$  multilayer stacks were also carried out. The structures of these multilayers (thicknesses, number of Co/Pt periods) were chosen around the values normally used for magnetooptic recording media. Compared with the simple picture of the variation in anisotropy with Co thickness, explained above in the case of the simple films, the effects of the irradiation on the magnetic properties are made more complex in multilayers by the magnetic interaction between the layers, which may be bipolar in origin, or an exchange interaction carried by the conduction electrons in the platinum. The latter interaction, which is actually manifested by ferromagnetism of the Pt for the interface layers, helps to raise the Curie temperature of the multilayers, especially when the Co thickness is very small. The presence of these two interactions also leads to the existence of quite a wide Co thickness range in which the system is decomposed into regular magnetic domains within which the magnetization is perpendicular ("strip" domain configuration), even for

slightly negative  $K_{eff}$  values where an easy magnetization plane would be expected.

The tests were carried out on two series of specimens, of the same Co thickness (and therefore the 5 same single layer anisotropy) and the same number of periods, but differing in the thickness of the Pt separating layer:

A series: Pt(2 nm)/[Pt(1.4 nm)/Co(0.3 nm)]<sub>6</sub>/Pt(6.5 nm)

B series: Pt(2 nm)/[Pt(0.6 nm)/Co(0.3 nm)]<sub>6</sub>/Pt(6.5 nm)

10 In the case of the B series, the Pt concentration of the alloy after complete interdiffusion would be about 66% (ferromagnetic alloy) while it would be 82% for the A series (nonmagnetic alloy). On the other hand, in the B series, in which 15 the Pt interlayer is thinner, the Co layers are more highly interacting, which in principle makes it easier to obtain the "strip" domain configuration, followed by the easy magnetization plane, by a reduction in the anisotropy.

20 Over the range of doses tested (up to  $10^{16}$  in the case of the A series and  $2.6 \times 10^{16}$  in the case of the B series), the irradiation results show qualitatively the same effects for both series: gradual (and easily controllable) transition from a perpendicular easy 25 magnetization axis (with a perfectly square hysteresis loop whose coercive field decreases with the irradiation dose) to a "strip" domain configuration, and then to an easy magnetization plane. As explained above, this tilting takes place at a lower dose for the 30 B series ( $3 \times 10^{15}$  as opposed to  $6 \times 10^{15}$  ions/cm<sup>2</sup>). At the doses used, all the specimens remained ferromagnetic at room temperature.

In all the cases described above, no variation in the surface roughness of the specimen could be 35 detected by AFM in air, even for extremely low, of the order of 0.2 nm rms, initial roughnesses.

Tests with irradiation through a resin mask were also carried out.

On Pt (3.4 nm) / Co (0.5 nm) / Pt (6.5 nm) / Herasil simple sandwich specimens, two types of resin were tested:

1. A Shipley negative resin, suitable for  
5 submicron lithography by X-ray lithography. The resin had been deposited as a thick (0.8  $\mu\text{m}$ ) layer over only half of a specimen and then annealed under the usual conditions. The entire specimen was then irradiated and the resin removed, again under the usual conditions  
10 (hot trichloroethylene bath).

The part unprotected by the resin reproduces the effects of the irradiation that were described above, whereas the protected part shows no change in its properties. In principle, using processes already  
15 developed elsewhere, the use of the same resin, but with in addition an X-ray lithography step in order to define an array of holes therein, should at the very least make it possible to obtain arrays of magnetically etched bits 0.2  $\mu\text{m}$  in size separated by 0.2  $\mu\text{m}$ , i.e. a  
20 recording density of 25 bits per  $\mu\text{m}^2$ , almost 20 times greater than the current densities;

2. a PMMA positive resin suitable for electron lithography. The resin was deposited as a layer about 0.85  $\mu\text{m}$  in thickness and in this case was not annealed,  
25 something which might have an influence on the quality of the pattern edges. Under the standard annealing conditions for this resin (160°C, 30 min) effects start to appear in the specimens, but annealing of just as good quality is possible at lower temperatures  
30 (<120°C), at which the specimens are insensitive). Next, the specimens underwent an electron lithography step in order to define, as recesses in the resin, an array of lines 1  $\mu\text{m}$  in width, separated by 1  $\mu\text{m}$ , over an area of 800×800  $\mu\text{m}^2$ . The entire specimen was then  
35 irradiated and the resin removed under the standard conditions. Observation in a magnetooptic microscope shows that, at the chosen irradiation dose ( $10^{16}$  atoms/cm $^2$ ), the irradiated part becomes paramagnetic at room temperature (this state has the

advantage of eliminating the coupling between magnetic regions). The part protected by the resin remains magnetized perpendicularly, with a square loop similar to that of the original specimen.

5        The same electron lithography process as above was applied to a Pt(2 nm)/[Pt(0.6 nm)/Co(0.3 nm)]<sub>6</sub>/Pt(6.5 nm) multilayer of the B series in order to create the same array of lines, followed by an irradiation at a dose of  $2 \times 10^{15}$  atoms/cm<sup>2</sup>. However, 10      unlike in the case of the single 0.5 nm Co layer, the two parts (the protected part and the irradiated part) retain a perpendicular magnetization and a square loop with, however, a lower coercive field in the case of the irradiated part. In fact, observation in a 15      magnetooptic microscope clearly shows a reversal of the magnetization in the reverse applied field after saturation, which firstly takes place in the irradiated lines and then propagates into the unirradiated parts (lines and film outside the array). In the intermediate 20      region, magnetic domains artificially created by lithography are therefore obtained. Next, tests were carried out using near-field magnetooptic microscopy, which made it possible to see these artificial domains very precisely. This consequently demonstrates the 25      feasibility of the proposed "contact" recording process. On the other hand, on specimens that were similar but were etched by material ablation, the same near-field microscopy technique reveals only the diffraction effects.

30       It should be noted that, after irradiation, the PMMA resin becomes more difficult to remove. Residues remaining along the features introduce roughness and a weak optical contrast of nonmagnetic origin, something which requires an additional stripping procedure in an 35      "oxygen plasma" (a procedure well known in microtechnologies).

Finally, given the precision of PMMA-resin electron lithography, we might expect to achieve bit

sizes of less than 100 nm, i.e. a density greater than 100 bits/ $\mu\text{m}^2$ .

The techniques of the type that have just been described are advantageously used for manufacturing 5 films which include buried magnetic structures, especially for the production of magnetically structured recording media or of magnetoelectronic devices, such as M-RAM memories, logic devices, etc.

They allow planar magnetic etching of buried 10 magnetic layers, which does not modify the surface roughness of the material and makes it possible to control the variations in optical properties, for example to make them negligible.

These techniques can be used for mass 15 production on an industrial scale.

Using light ions, which have no etching effect, these can be deeply implanted into the substrate, well below the layer.

The parameter is then the energy deposited per 20 ion along the trajectory - and not the cascades of defects generated by heavy ions - thereby allowing excellent control of the electromagnetic modifications, for high doses, something which gives a homogeneous effect.

25 Moreover, an easy nucleation region, due to the reversal of the magnetization) and associated with phenomena occurring at the border of the irradiated region, is intrinsically obtained with the proposed technique. This is a major advantage for controlling 30 and standardizing the magnetization reversal field in an assembly of magnetic "particles", either for a recording medium material or for a memory or logic chip, without limitation.